

# Independent Pair Parton Interactions Model and Multiplicity Distributions at LHC

I.M. Dremin, V.A. Nechitailo

Lebedev Physical Institute, Moscow



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The independent pair parton interactions (IPPI) model:

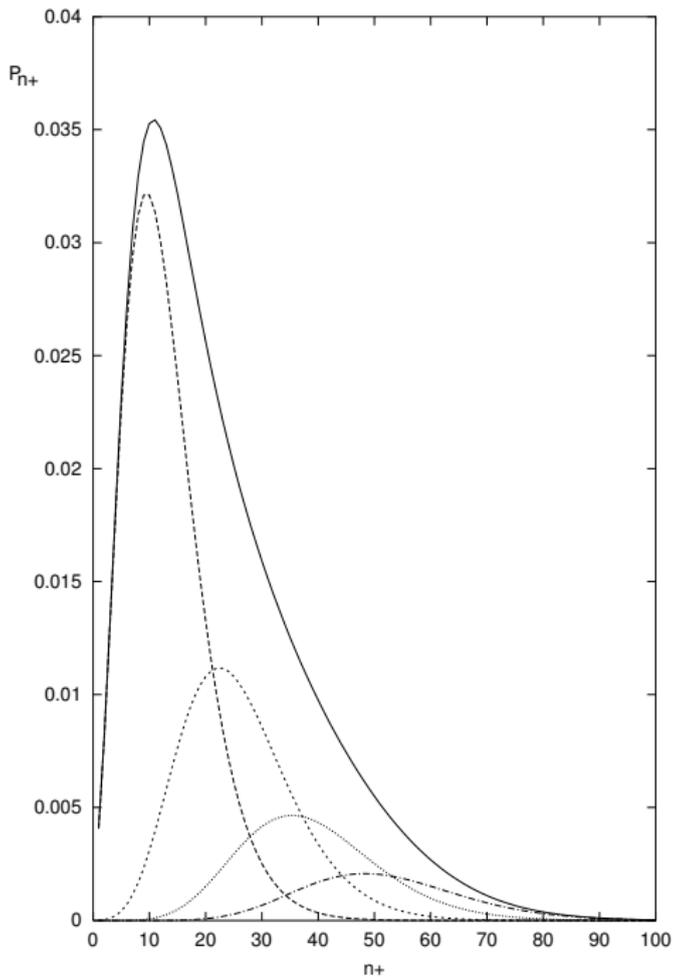
1. High energy particles are "clouds" of partons.
2. Each pair of colliding partons is independent of others and creates new particles according to NBD-distribution of multiplicity.

The main equation:

$$P(n; m, k) = \sum_{j=1}^{j_{max}} w_j P_{NBD}(n; jm, jk). \quad (1)$$

$P(n; m, k)$  is the probability to create  $n$  particles,  $m$  and  $k$  are the parameters of the NBD-distribution,  $w_j$  is the probability for  $j$  parton pairs to be active at a given energy,  $\sum_{j=1}^{j_{max}} w_j = 1$ ,  $j_{max}$  is the maximum number of the active parton pairs. The simplest case:  $w_j = w_1^j$ .

- Experimental indications: single NBD fits at energies up to 200 GeV, then the distributions widen.
- Interpretation: 1 pair of partons is active and leads to NBD at lower energies while their number increases with energy.



The decomposition of the multiplicity distribution at Tevatron  $\sqrt{s}=1.8$  TeV into 1, 2, 3 and 4 parton-parton interactions

How one gets the main equation.

Key property: Convolution of NBDs is again NBD!

$$P(n; m, k) = \sum_{j=1}^{j_{max}} w_j P_j(n; m, k) = \sum_{j=1}^{j_{max}} w_j \sum_{(n_p)} \prod_{p=1}^j P_{NBD}(n_p; m, k). \quad (2)$$

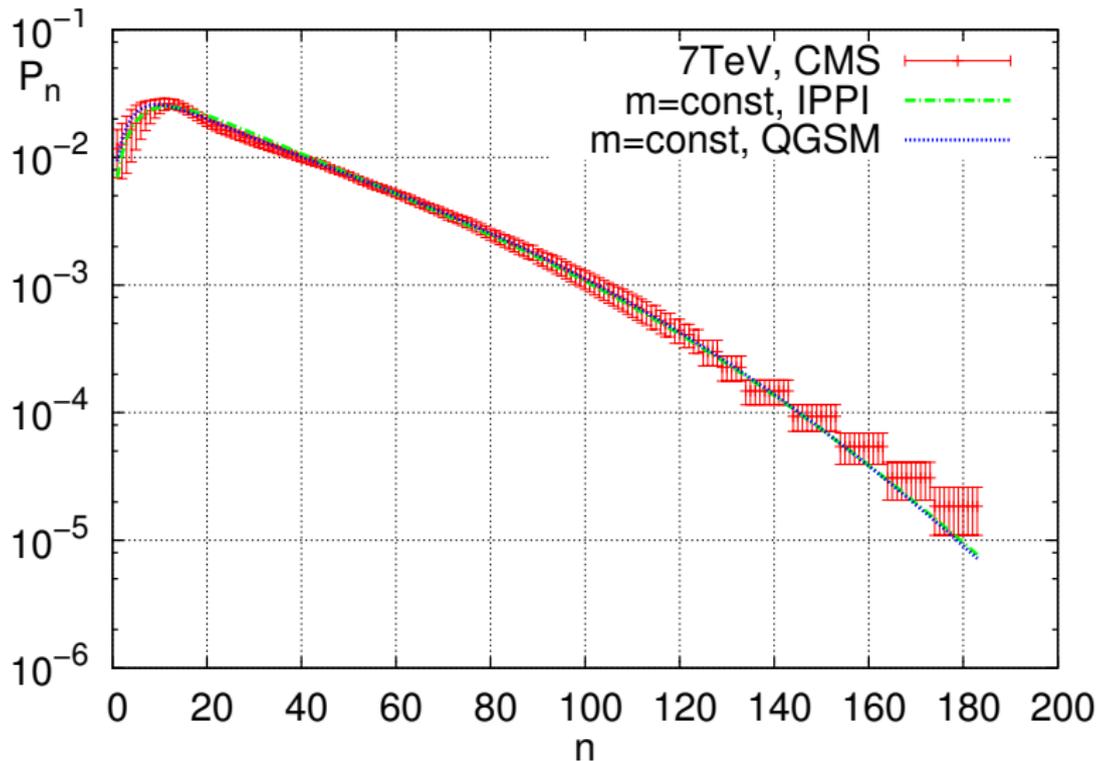
$n_p$  is the number of particles created by the  $p$ th pair,

$\sum_{(n_p)}$  denotes the convolution of NBD distributions with the sum of those parton interactions where  $\sum_{p=1}^j n_p = n$ .

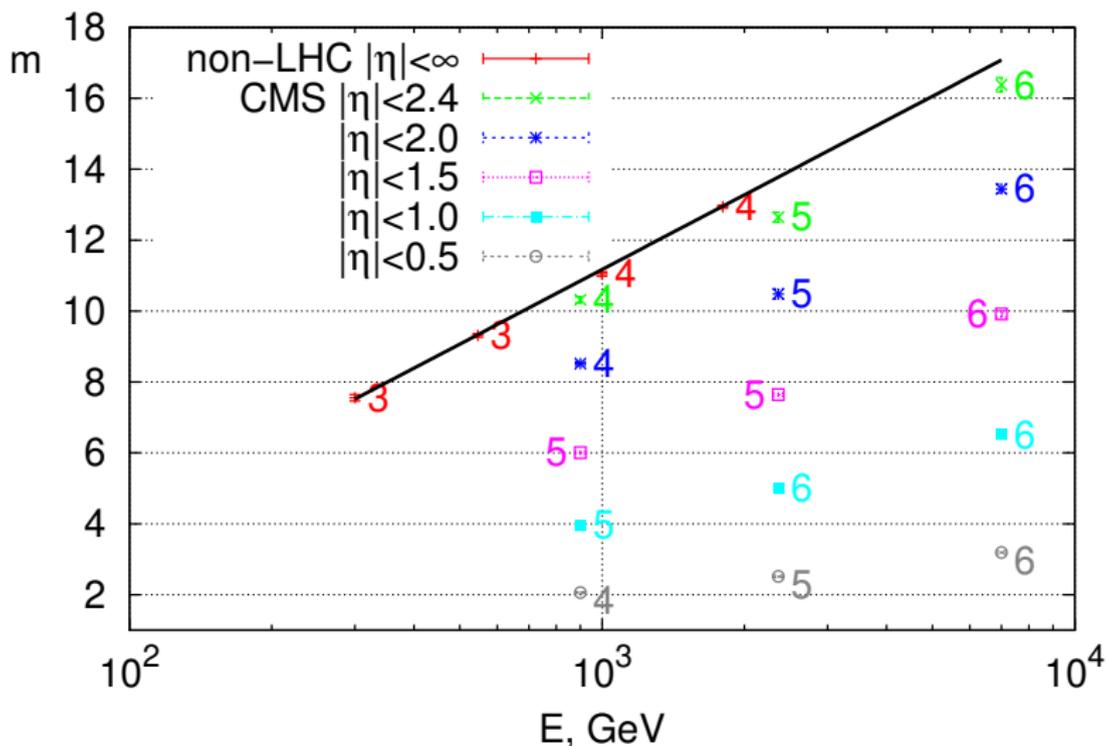
The summation in Eq. (2) gives rise to the main equation.

NBD definition:

$$P_{NBD}(n; m, k) = \frac{\Gamma(n+k)}{\Gamma(n+1)\Gamma(k)} \left(\frac{m}{k}\right)^n \left(1 + \frac{m}{k}\right)^{-n-k}$$



The fits by the **IPPI**-model (line) and by the **QGSM**-model (dots) of the multiplicity distribution at 7 TeV ( $|\eta| < 2.4$ , CMS-data).



The values of the parameter  $m$  ( i.e. of the effective average multiplicity for a single parton pair ) at different energies and rapidity windows.

The number of active pairs ( $j_{max}$ ) is shown near each point.

## IMPORTANT TECHNICALITIES!

The requirement of independence of  $m$  on the ranks of moments of the distribution imposes restrictions on the parameter  $k$ .

We use the properties of factorial ( $F_q$ ) and cumulant ( $K_q$ ) moments of the multiplicity distributions (as well as of their ratio  $H_q$ ) in fits of experimental data to show how well this requirement is fulfilled.

For more details see the papers in Phys. Rev. or arXiv:hep-ph.

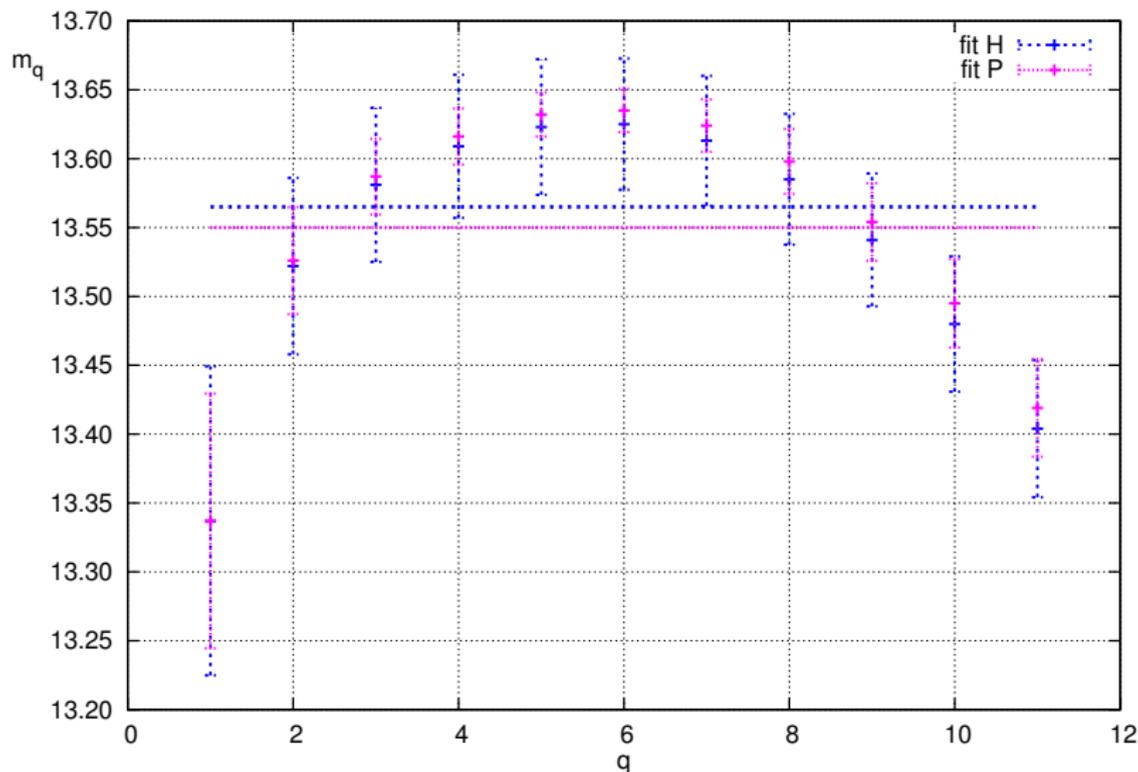
$$F_q = \sum_n P(n)n(n-1)\dots(n-q+1)$$
$$= \sum_{j=1}^{j_{\max}} w_j \frac{\Gamma(jk+q)}{\Gamma(jk)} \left(\frac{m}{k}\right)^q = f_q(k) \left(\frac{m}{k}\right)^q; \quad K_q = \kappa_q(k) \left(\frac{m}{k}\right)^q \quad (3)$$

where

$$f_q(k) = \sum_{j=1}^{j_{\max}} w_j \frac{\Gamma(jk+q)}{\Gamma(jk)} = k \sum_{j=1}^{j_{\max}} w_j j(jk+1)\dots(jk+q-1). \quad (4)$$

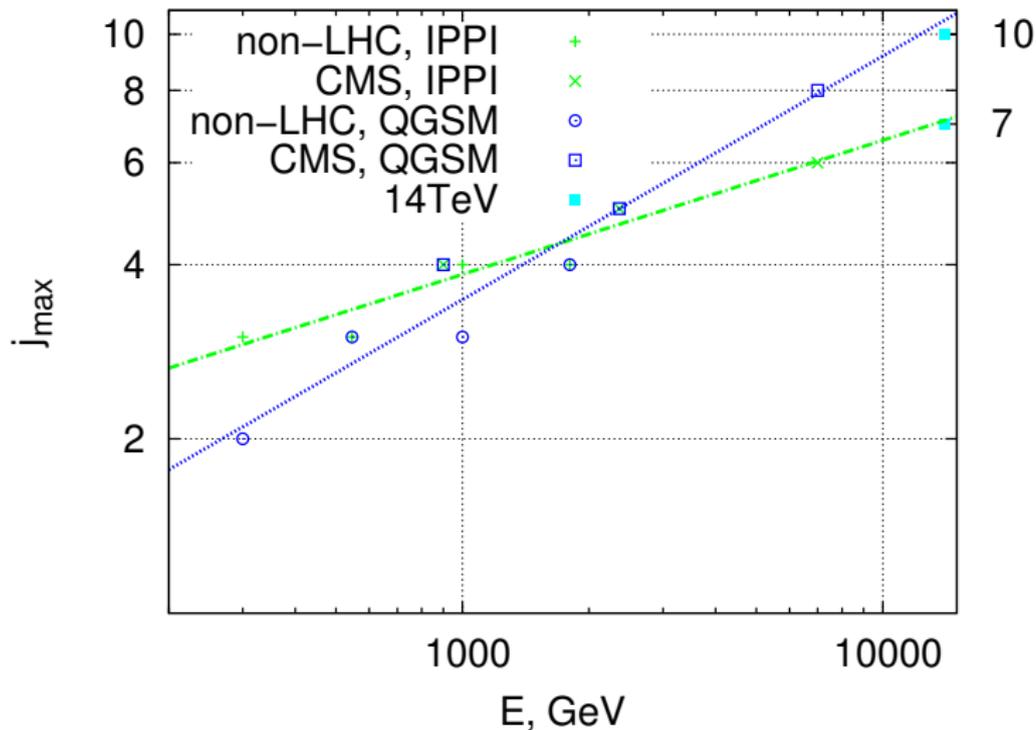
The selfconsistency of the IPPI-model asks for  $q$ -independence of  $m$ :

$$m = k \left( \frac{F_q^{exp}}{f_q(k)} \right)^{1/q} = \text{const.} \quad (5)$$



- The IPPI-model is proposed.
- The pairs of partons from colliding "clouds" are independent and each pair creates particles according to NBD-distribution.
- Experimental distributions are well described at different energies and rapidity windows with only two adjustable parameters  $k$  and  $j_{max}$ .
- The average multiplicity in collision of a single pair  $m$  and the number of active pairs  $j_{max}$  increase with energy logarithmically.
- The density of the parton medium increases with energy and asks for account of SOFT (not only HARD) multiparton interactions.
- The predictions at higher energies 14 and 100 TeV have been done.

ADDITIONAL SLIDES



Extrapolation of the maximum number of the active parton pairs ( $j_{max}$ ) to 14 TeV.

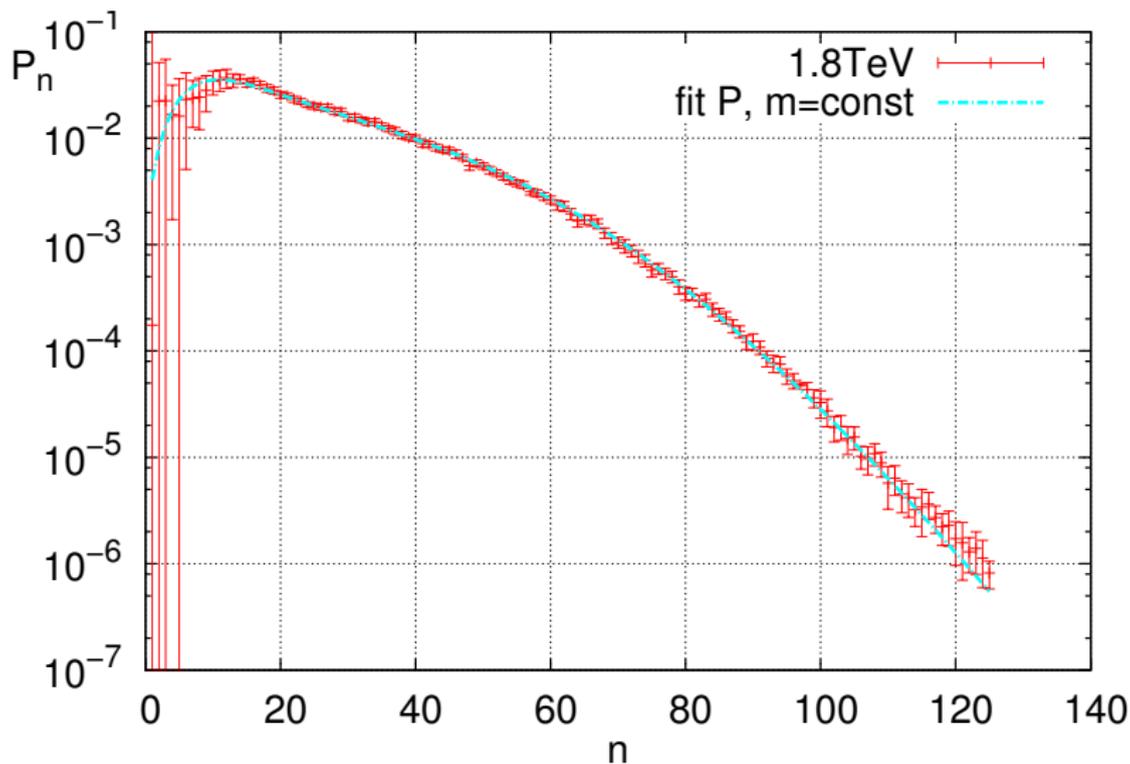
QGSM (A.B. Kaidalov, K.A. Ter-Martirosyan)

$$w_j(\xi_j) = \frac{p_j}{\sum_{j=1}^{j_{\max}} p_j} = \frac{1}{jZ_j(\sum_{j=1}^{j_{\max}} p_j)} \left( 1 - e^{-Z_j} \sum_{i=0}^{j-1} \frac{Z_j^i}{i!} \right) \quad (6)$$

where

$$\xi_j = \ln(s/s_0 j^2), \quad Z_j = \frac{2C\gamma}{R^2 + \alpha'_P \xi_j} \left( \frac{s}{s_0 j^2} \right)^\Delta \quad (7)$$

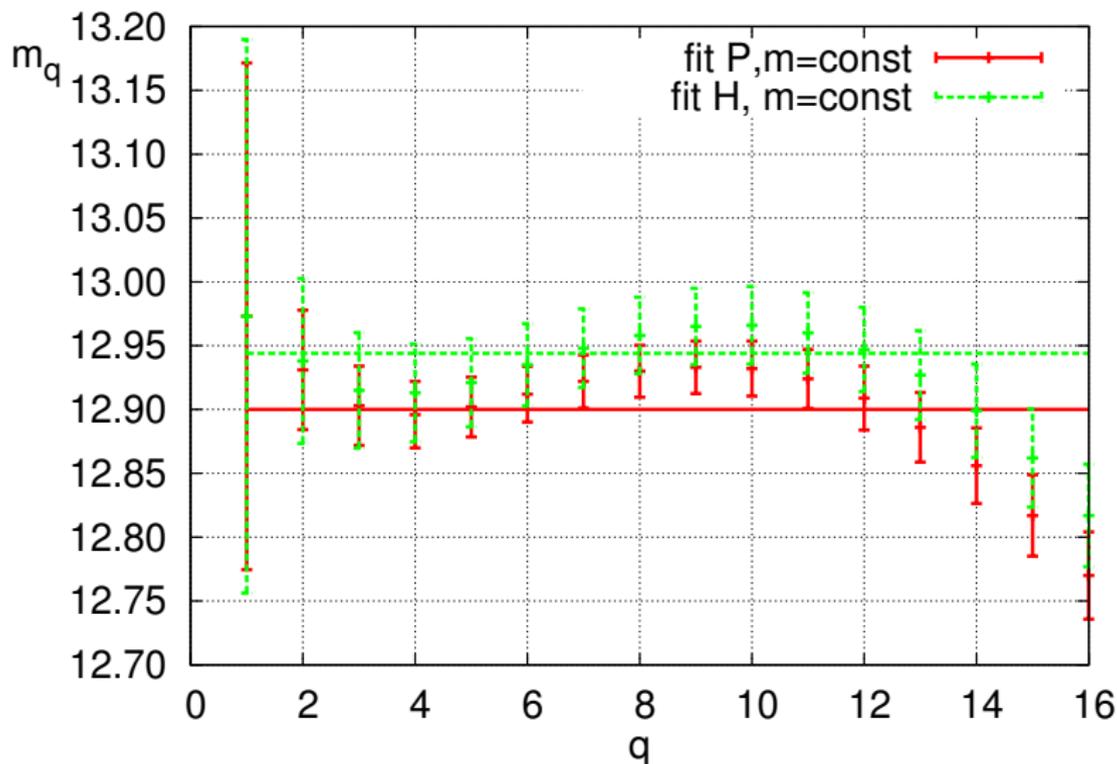
numerical parameters derived from fit of experimental data on total and elastic scattering  $\gamma = 3.64 \text{ GeV}^{-2}$ ,  $R^2 = 3.56 \text{ GeV}^{-2}$ ,  $C = 1.5$ ,  $\Delta = \alpha_P - 1 = 0.08$ ,  $\alpha'_P = 0.25 \text{ GeV}^{-2}$ ,  $s_0 = 1 \text{ GeV}^2$ .



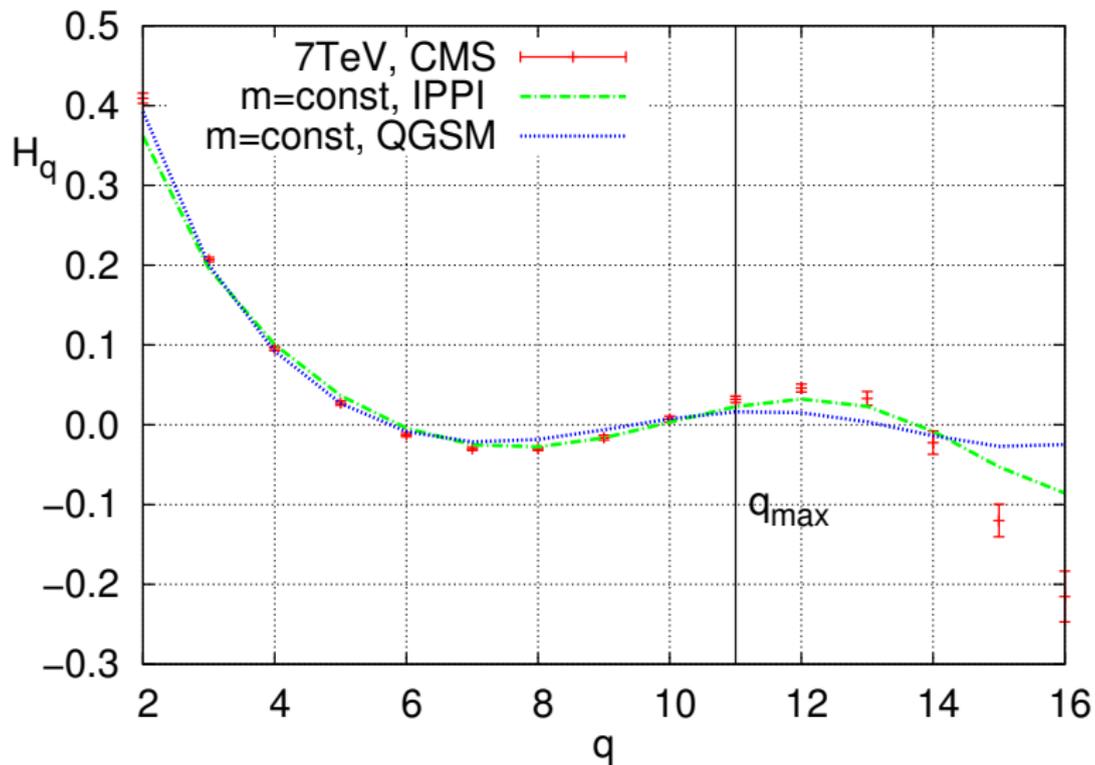
The fit by the IPPI-model (dotted) of the multiplicity distribution at 1.8 TeV.

The selfconsistency of the IPPI-model asks for  $q$ -independence of  $m$  (1.8 TeV):

$$m = k \left( \frac{F_q^{\text{exp}}}{f_q(k)} \right)^{1/q} = \text{const.} \quad (8)$$



$$K_q = F_q - \sum_{r=1}^{q-1} \frac{(q-1)!}{r!(q-r-1)!} K_{q-r} F_r. \quad (9)$$



IPPI-fit (dash-dotted) and QGSM-fit (dotted) for  $H_q$ -moments at 7 TeV.